Memory Formation in Driven Disordered Systems – Dead or Alive

Muhittin Mungan

Institute of Biological Physics U. Cologne

Les Houches 2023

M. Mungan ()

Memory Formation in Matter

 Suppose we load a spring and a sand bag with varying loads ...





• Place two identical medium loads on each ...





- Remove loads.
- Spring restores initial position, sand bag remains deformed.





- Next place identical lighter loads on spring and bag.
- Spring compresses less, while sand bag retains deformed state.





- Removing the lighter load, spring returns to initial position.
- Sand bag retains initial deformation:
- Sand bag keeps a **memory** of initial loading.

-	
_	
-	
_	
_	
-	
_	
-	
_	
_	



- Place next identical heavier loads.
- Spring compresses more than before.
- Sand bag also deforms further.





- Remove heavy loads.
- Spring **restores** its uncompressed state.
- Sand bag remains in further deformed state.



Sand bag:

- Retains memory of past deformations.
- Placing heavier load on sand bag, ⇒ memory of previous load is lost.
- Memory is overwritten by heavier load.



Two "devices", two different responses:

• Spring: perfectly elastic: returns to same position when unloaded. No memory of its loading history.





- **Spring:** "records" instantaneous state of loading, *e.g.* like a kitchen scale.
- Sand bag: "records" extreme loading event in its history ⇒ future response is history-dependent.

- Example of system that can record the largest applied load.
- Simple prototype of a system interacting with a changing environment:
 - System: sand bag,
 - Changing Environment: various loads placed on sand bag,
 - Effect on system: altering shape of sand bag.
- Main ingredients:
 - Disorder
 - Large number of degrees of freedom

Q: How do we characterize systems that retain a memory of their past environments?

- Memory formation in driven soft-matter systems.
- Focus on athermal, quasi-static (AQS) response to driving.
- Response to AQS driving can be captured via state transition graphs: ⇒ Dynamical features are encoded in graph topology
- Demonstrate how these ideas can be used to **understand and utilize** memory formation.
- Applications: Show that these ideas can be used to analyze
 - dynamics of a sheared amorphous solid,
 - biological evolution in changing environments.

Origami-bellows as a mechanical memory device









Preisach Element (Hysteron)



(Jules, Reed, Daniels, MM, & Lechenault Phys. Rev. Res. (2022))

A stack of four bellows - The Preisach Model



(Jules, Reed, Daniels, $\mathbf{MM},$ & Lechenault Phys. Rev. Res. (2022))

Les Houches 2023

∃ ▶ ∢

A stack of four bellows - The Preisach Model



(Jules, Reed, Daniels, MM, & Lechenault Phys. Rev. Res. (2022))

Image: Image:

Les Houches 2023

3 1 4





Single site flips suffice to regain stability NO AVALANCHES!

(M.M. Terzi & **MM** PRE **102** (2021) 012122)



(M.M. Terzi & MM PRE 102 (2021) 012122)

Suppose we have a system with N hysterons: $\boldsymbol{\sigma} = (\sigma_1, \sigma_2, \dots, \sigma_N)$ is a state. $\sigma_i = 0.1$ Switching fields: $F_1^{\pm}, F_2^{\pm}, \ldots, F_N^{\pm}$ F^+ $F_{i}^{-} < F_{i}^{+}$ $\sigma_i = 1 \text{ requires } F > F_i^ \sigma_j = 0 \text{ requires } F < F_j^+$ POSSIBLE ONLY IF σ such that: $F^{-}[\boldsymbol{\sigma}] \equiv \max_{\substack{\{i : \sigma_i = 1\}}} F_i^{-} < \min_{\substack{\{j : \sigma_j = 0\}}} F_j^{+} \equiv F^{+}[\boldsymbol{\sigma}]$ (Stability condition, determines the set of states) By INDEPENDENCE label hysterons s.t. Up: $1 \rightarrow 2 \rightarrow \ldots \rightarrow N$ Down: $\rho_1 \to \rho_2 \to \ldots \to \rho_N$ $\mathbf{F}_{\partial \mathbf{N}}^{-} \cdots \mathbf{F}_{\partial \mathbf{n}}^{-} \mathbf{F}_{\partial \mathbf{n}}^{-} \mathbf{F}_{\mathbf{n}}^{+} \mathbf{F}_{\mathbf{n}}^{+} \mathbf{F}_{\mathbf{n}}^{+} \mathbf{F}_{\mathbf{N}}^{+} \mathbf{F}_{\mathbf{n}}^{+}$ Switching Sequence ρ determines transition graph between (00...0) and (11...1)!(M.M. Terzi & **MM** PRE **102** (2021) 012122)

M. Mungan ()



M. Mungan ()



M. Mungan ()



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$

Switching Sequence ρ COMPLETELY determines the transition graph between (00...0) and (11...1)!



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$

loop Return Point Memory (ℓ RPM) as topological feature of transition graph

⇒ Hierarchical Structure of loops nested within loops "Every loop is existed from its end points!"



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$

Given ρ what is the number of vertices in the main loop?

M.M. Terzi & **MM** PRE **102** (2021) 012122,
P. L. Ferrari, **MM** & M.M. Terzi AIHPD 2022

Memory Formation in Matter



Switching Sequence:

Up: $1 \rightarrow 2 \rightarrow 3 \rightarrow 4 \rightarrow 5 \rightarrow 6 \rightarrow 7 \rightarrow 8$

Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$

Given ρ what is the number of vertices in the main loop? ANSWER: It is equal to the # of increasing subsequences contained in ρ

M.M. Terzi & **MM** PRE **102** (2021) 012122,
P. L. Ferrari, **MM** & M.M. Terzi AIHPD 2022



Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$

Given ρ what is the number of vertices in the main loop? ANSWER: It is equal to the # of increasing subsequences contained in ρ MOREOVER: each increasing subsequence encodes a deformation history!

M.M. Terzi & **MM** PRE **102** (2021) 012122,
P. L. Ferrari, **MM** & M.M. Terzi AIHPD 2022

Memory Formation in Matter



Down: $2 \rightarrow 1 \rightarrow 8 \rightarrow 4 \rightarrow 5 \rightarrow 3 \rightarrow 6 \rightarrow 7$

Given ρ what is the number of vertices in the main loop? ANSWER: It is equal to the # of increasing subsequences contained in ρ MOREOVER: each increasing subsequence encodes a deformation history!

Histories are mapped into states

M.M. Terzi & **MM** PRE **102** (2021) 012122, P. L. Ferrari, **MM** & M.M. Terzi AIHPD 2022

Recap: Memory formation in driven disordered systems

- Return point memory (RPM) is one way of encoding memory of deformation history.
- Memory and history-dependence become topological features of transition graph (*t*-graph) ⇒ loop return point memory (*ℓ*RPM), MM & Terzi, AHP (2019).
- Preisach model: **simplest model with** *l***RPM**.
- More complicated *l*RPM systems: Random field Ising model with ferromagnetic interactions, models of depinning, i.e. elastic manifolds in random media. Consequence of no-passing (Middleton 1992): partial order on configurations that is preserved by dynamics.
- Systems that exhibit *l*RPM approximately: the sheared amorphous solids (MM, S. Sastry, K. Dahmen & I. Regev, PRL (2019)).

Where to go from here

• The Preisach Model:

- Hystorically, introduced to describe hysteresis in magnetic materials (Preisach 1935).
- "Ideal gas" of memory formation:
 - Hysterons do not interact, switching thresholds do not depend on states of the other hysterons. \Rightarrow switching order ρ fixed.
- Hysteron interactions ⇒ rich & versatile dynamics: hysterons can be created, destroyed and altered.
- Possible directions:
 - Device Design: mechanical sensors, automata & A/D converters of environmental signals, control systems (Lechenault & van Hecke groups),
 - Adaptive evolution in a changing environment: (Das, Krug & MM, 2022)

• Understanding the dynamics of **non-quenched**, i.e. **co-evolving** (annealed) disorder in driven disordered systems: yielding and irreversibility transition in sheared amorphous solids.

Where to go from here? Understanding annealed disorder

Periodically compressing/uncompressing a crumpled sheat:

(Shohat et al. PNAS 2022)



Cyclic Response achieved after a number of driving cycles
Where to go from here? Understanding annealed disorder

Periodically compressing/uncompressing a crumpled sheat:

(Shohat et al. PNAS 2022)



Cyclic Response achieved after a number of driving cycles

Cyclic Response has discontinuities Can be traced back to creases and wrinkles: localized, bi-stable elements \Rightarrow hysterons

Mechanical annealing:

- \Rightarrow interacting hysteron system
- \Rightarrow hierarchy of nested cycles

Where to go from here? Understanding annealed disorder

Periodically compressing/uncompressing a crumpled sheat:

(Shohat et al. PNAS 2022)



Cyclic Response achieved after a number of driving cycles

Cyclic Response has discontinuities Can be traced back to creases and wrinkles: localized, bi-stable elements \Rightarrow hysterons

- Mechanical annealing:
 - \Rightarrow interacting hysteron system
 - \Rightarrow hierarchy of nested cycles



Where to go from here? Understanding annealed disorder

Periodically compressing/uncompressing a crumpled sheat:

(Shohat et al. PNAS 2022)



Cyclic Response achieved after a number of driving cycles

Cyclic Response has discontinuities Can be traced back to creases and wrinkles: localized, bi-stable elements

 \Rightarrow hysterons

Mechanical annealing:

- \Rightarrow interacting hysteron system
- \Rightarrow hierarchy of nested cycles

Simulations: (Kumar et al. JCP 2022) Periodically shearing an amorphous solid



- Under cylic loading \Rightarrow evolution into periodic response.
- Periodic response \Rightarrow sequence of **mechanical instabilities**.
- Periodicity: these instabilities are repeatedly retriggered, ⇒ reversible instabilities.
- Transient towards periodicity: Trial & Error or Evolution?

What characterizes the attainment of periodic response?

What characterizes the yielding transition?

Probing molecular glasses via t-graphs

- Focus on amorphous solids under externally applied shear strain:
 - Amorphous = non-crystalline structure, **disordered**.
 - Numerically achieved by considering binary mixtures of particles of **different sizes**, prevents crystallization.
 - Particles interact with each other.
- AQS dynamics: (Maloney & Lemaître PRE 74 (2006) 016118):
 - Athermal: ignore thermal activation processes ⇒ mechanical equilibria
 - Slowly varying shear strain γ_t: fast relaxation to new mechanically stable configuration ⇒ rate-independent response ⇒ quasi-static process.



The Yielding transition



The Irreversibility transition

Cyclic Response under Oscillatory Shear Strain:



Irreversibility Transition: Transient to limit-cycle diverges as $\gamma_{\text{max}} \nearrow \gamma_{\text{c}}$.

Pine, Gollub, Brady & Leshansky Nature **438** 997 2005; Corte, Chaikin, Gollub & Pine Nat. Phys. **4** 420, 2008 ...

How can we understand the evolution?

t-graphs and mesostates

Configuration Space for N-particles: Γ N-particle configuration: *x* Deformation parameter: γ

AQS dynamics:

 $E = U(\boldsymbol{x}, \boldsymbol{\gamma})$

(potential energy)

Stable Configurations x evolve with γ until they become unstable:

 $\Rightarrow \boldsymbol{x}(\gamma) \text{ for } \gamma^- < \gamma < \gamma^+$



plastic events are localized rearrangements:

Mesostate transitions are plastic events!



 $A \rightarrow B$







M. Mungan ()

Memory Formation in Matter

Les Houches 2023

Image: A matrix and a matrix



M. Mungan ()

Memory Formation in Matter

M. Mungan ()

Memory Formation in Matter

And a decision of the second s



M. Mungan ()

Memory Formation in Matter



M. Mungan ()

Memory Formation in Matter







M. Mungan ()

Memory Formation in Matter



The lay of the land — g = 25, $\mathcal{N} = 1416$ mesostates



- Degeneracy: many deformation paths lead to the same state ${\cal N}=1416\ll 2^{26}\approx 6\times 10^7$
- Bottlenecks, tree-like regions & loops

Transients and limit-cycles





Oscillatory Shear does not just lead to cyclic response! We get more than we asked for: a hierarchy of nested cycles

Mungan et al. PRL 2019



Oscillatory Shear does not just lead to cyclic response! We get more than we asked for: a hierarchy of nested cycles

Periodically compressing/uncompressing a crumpled sheet Cyclic Response achieved after a number of driving cycles (Shohat et al. PNAS 2022)



Mungan et al. PRL 2019

M. Mungan ()



Oscillatory Shear does not just lead to cyclic response! We get more than we asked for: a hierarchy of nested cycles



Mungan et al. PRL 2019



Oscillatory Shear does not just lead to cyclic response! We get more than we asked for: a hierarchy of nested cycles



Mungan et al. PRL 2019





Oscillatory Shear does not just lead to cyclic response! We get more than we asked for: a hierarchy of nested cycles



M. Mungan ()



- Mechanical Aging: Emergence of interacting soft-spot system
- Persistence & Robustness: Soft-spots play nice (most of the time)
 ⇒ hierarchy of nested cyles
- Selection & Memory: Emergent soft-spot system determined by form of periodic loading, e.g. amplitude of oscillation ...

Down the rabbit hole – total irreversibility



- A pair of mesostates, A and B, is **mutually reachable**, if there exists some deformation path from A to B AND from B to A
- Mutual reachability (MR) is an equivalence relation \Rightarrow partition of ${\cal S}$
- Equivalence classes under MR are called **strongly connected components (SCCs)**
- Toplogy \leftrightarrow Physics:
 - Transitions within an SCC are reversible, (reversible plasticity)
 - Transitions between an SCC are irreversible, (irreversible plasticity)
 - Any periodic response must be confined to a single SCC

SCCs, reversible and irreversible transitions



- Irreversibility encoded in the transitions between SCCs
- \Rightarrow inter-SCC transition graph (acylic)

SCC size distribution and length of transients



- cyclic response must be confined to a single SCC
- large amplitude cyclic response requires large SCCs
- large SCCs are rare, especially close to yielding (Regev et al. 2021)
- transient length increase as amplitude increases

A fitness landscape model describing the evolution of antibiotic drug resistance:

- The trade-off induced fitness landscape model (TIL) (S.G. Das, S.O.L. Direito, B. Waclaw, R. Allen & J. Krug, eLife 9 (2020) e55155):
 - Bacteria in environment of varying antibiotic drug concentration x.
 - L possible loci where mutations can occur.
 - Binary vector σ = (σ₁,..., σ_L) encodes whether mutation at *i* is present (σ_i = 1) or absent (σ_i = 0).
- Environment characterized by **single parameter:** antibiotic concentration *x*.
- For each x: a mapping that assigns to each genotype σ a fitness f_{σ} .

A fitness landscape model describing the evolution of antibiotic drug resistance

- **Goal:** Characterize the transition between genotypes as *x* is changed and fitness maxima change.
- Motivation:



- L = 4 mutation sites associated with an antibiotic resistance enzyme when subject to antibiotic at different concentrations.
- Black/Gray and Red/Orange arrows indicate transitions to new fitness peaks under incease and decease of concentration.
- Data compiled from M. Mira *et al.* Mol. Bio. Evol. **32** (2015) 2707.

The Das-Krug *et al.* trade-off induced fitness landscape model (TIL) as a system of interacting hysterons.



∃ ▶ ∢ ∃

TIL nested cycles



M. Mungan ()

Memory Formation in Matter
Outlook & Conclusions

- Overview of memory formation in driven soft-matter systems.
- Approach via state transition graphs.
- Glassy systems like the sheared amorphous solids:
 - Hysteron interactions \Rightarrow rich & versatile dynamics
 - The interacting hysteron system is often emergent: The exposure to external driving or interactions with environment shapes the formation of hysteron systems ⇒ Aging
 - Quenched vs. Annealed disorder.
- Parallels with adaptive evolution in biology.
- Connection with biology goes both ways, e.g. understanding evolution of antibiotic resistance:

PHYSICAL REVIEW X 12, 031040 (2022)

Featured in Physics

Driven Disordered Systems Approach to Biological Evolution in Changing Environments

Suman G. Dæ,¹ Joæchim Knug⁰,¹ and Muhitiin Mungan^{0,12} ¹Institute for Biological Physics, University of Cologne, Zülpicher Strafe 77, D-50937 Köln, Germany ²Institut für Angewandte Mathematik, Universitä Bonn, Endersicher Allee 60, D-53115 Bonn, Germany

THANK YOU!

• Collaborators:

- RPM & Preisach Model: M. Mert Terzi (ESPCI)
- Origami: Theo Jules, Frederic Lechenault (ENS Paris), Austin Reed (Indiana U.), K. Daniels (NCSU)
- Amorphous Solids Atomistic systems: Ido Regev, Ido Attiah, Asaf Szulc (Ben Gurion U.), Srikanth Sastry (JNCASR, India), Karin Dahmen (U. Illinois)
- Amorphous Solids Mesoscopic systems: Damien Vandembroucq, Dheeraj Kumar, Sylvain Patinet (ESPCI), Craig E. Maloney (Northeastern).
- Biological Systems: Joachim Krug, Suman Das (U. Cologne)