Cavitation instabilities in amorphous solids

Umang A. Dattani, IMSc, Chennai, India Interaction, Désordre, Elasticité (GDR-IDE), Les Houches, 3rd April 2023

In collaboration with: Rishabh Sharma, Smarajit Karmakar & Pinaki Chaudhuri

Introduction: amorphous solids

- Amorphous solids lack long-range order
- They exist in diverse forms





Introduction

Amorphous solids known to fracture via coalescence of smallcavities/voids



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SCIENCE ADVANCES | RESEARCH ARTICLE

CONDENSED MATTER PHYSICS

Observation of cavitation governing fracture in glasses

Lai-Quan Shen^{1,2,3}, Ji-Hao Yu^{1,3}, Xiao-Chang Tang^{1,3}, Bao-An Sun^{1,2}*, Yan-Hui Liu^{1,2,3}, Hai-Yang Bai^{1,2,3}*, Wei-Hua Wang^{1,2,3}

Background: Numerical simulations

- Intersection of glass-transition line with spinodal at finite temperature(Sastry, PRL, 1999)in the $T-\rho\,$ plane
- Mechanical instability below glass-transition temperatures due to gas-glass coexistence

Phase diagram of attractive glass-forming system



Testard, Berthier & Kob, PRL, 2011



 Cavitation instabilities in amorphous solids under uniform expansion

Umang A. Dattani, Smarajit Karmakar & Pinaki Chaudhuri. Phys. Rev. E 106, 055004, 2022.

 Cavitation instabilities in amorphous solids via secondary mechanical perturbations

Umang Dattani, Rishabh Sharma, Smarajit Karmakar & Pinaki Chaudhuri. arXiv: 2303.04529, March, 2023

Protocol: athermal quasistatic expansion

- Initial states prepared by cooling a high temperature liquid (MD, 2D Lennard-Jones system)
- Protocol: apply a constant volume strain in each step of the simulation followed by energy minimization(conjugate gradient)



Stress-strain curves

- On expansion pressure decreases, even goes negative
- A large pressure jump beyond a certain density
- Pressure jumps correspond to cavitation
- Cavities eventually merge and lead to complete fracture [(c)-(h)]





Pressure fluctuations

- Ensemble fluctuations in pressure as function of density define a susceptibility $\chi_p(\rho) = N \left(\langle P^2(\rho) \rangle - \langle P(\rho) \rangle^2 \right)$
- Sharp peak with increasing system sizes implies a yielding like transition
- The peak height goes as $\chi_p^{peak} \sim N$; brittle-like yielding with a macroscopic pressure drop (Ozawa et al, PNAS 2018)



Pressure jumps and plastic events

- Pressure jumps correspond to irreversible plastic events
- We diagonalize the hessian of potential energy to probe the stability near a plastic drop $H_{ij}^{\alpha\beta} = \frac{\partial^2 U}{\partial x_i^{\alpha} \partial x_j^{\beta}}$
- Near each plastic jump, the lowest non-zero eigenvalue of the hessian goes to zero as a power law: $\lambda_{min} \sim \sqrt{(\rho \rho_c)/\rho_c}$
- Saddle-node bifurcation on the potential energy landscape; seen under shear too



(Maloney & Lemaitre, PRE 2005)



eigenmodes & displacement fields

- On approach to a plastic jump, the eigenmode of λ_{min} predicts the displacement field
- Across the plastic jump, the displacement field has a smaller overlap with the eigenmode; cascade/avalanche nature of cavitation



Statistics of avalanches: post-yield

- Absence of a post-yield steady state
- Distributions of size of pressure jumps and energy drops collapse on scaling N by exponents -0.34 and 0.53
- The mean sizes of pressure jumps and energy drops too scale with same exponents



11

Pre-yield plastic events

- λ_{min} goes to zero multiple times even before cavities start to form
- These plastic events have a quadruopolar shape just like the events seen under shear $\rho_{1.12}$ $\rho_{1.15}$ $\rho_{1.18}$ 2 4 6





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Motivation

Negative pressure states in liquids are metastable(maxwell's construction), but our the liquid is below the glass-transition temperature?



 In real-life scenarios, the modes of deformation are always a combination of volume strains and shear strains, how to account for this?

Part II: Secondary defomation of the expanded solid

- Uniform expansion states subject to cyclic shear(T=0) and local random deformation via activity(T=0.01)
- Control parameters:
 - Amplitude of cyclic shear: γ_{max}
 - Magnitude of active force: f_0



Results: high density cavitation

Cavitation at higher densities for both secondary deformations



Energy barriers

- Lower energy barriers/thresholds to cavitation under secondary deformations
- Shear deformation and local random deformation couple better to cavitation instabilities than expansion



Yielding under cyclic shear

- Stroboscopic energy relaxation shows yielding under cyclic shear
- No limit cycles above yielding
- Cusp in average steady-plateau energy(Bhowmik, Foffi & Sastry, PNAS, 2021)
- Divergence of timescales to reach limit cycles with an exponent -2.8







Dependence on control parameters

- Cavitation region marked in two-parameter phase diagrams
- Cavitation occurs both, below and above, yield / fluidisation thresholds



Soft spots

 For a same expansion trajectory the soft spots for different phase parameters are different - abundance of soft cavitation modes



Spatial maps of plasticity: squared displacements

- Cyclic shear: system spanning avalanche-like structures at high values of γ_{max}
- Local random deformation: spatially localised bursts of plasticity
- Both local and systemspanning pathways lead to cavitation under secondary deformations





- Under expansion, the amorphous solids cavitate via a yielding-like transition
- The cavities grow with expansion resulting in a system-spanning crack
- Plastic pressure jumps correspond to saddle-node bifurcations on the PEL
- Early onset of cavitation for both secondary deformations, cyclic shear & local random deformation
- Lower energy barriers to cavitation under cyclic shear & local random deformation than expansion

